

# Determining the Nonperturbative Collins-Soper Kernel From TMD Observables Using Lattice QCD

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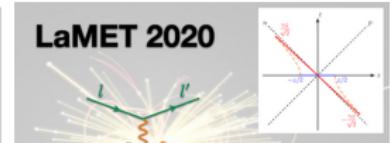
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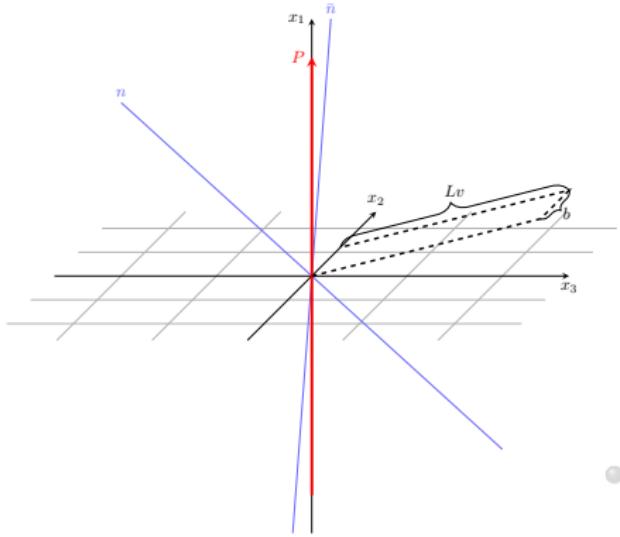
**LaMET 2020**



- 1 Lattice calculations
- 2 Fit correlators to parametrization to obtain the moments of TMDPDFs
- 3 Extraction of the CS-Kernel from lattice data

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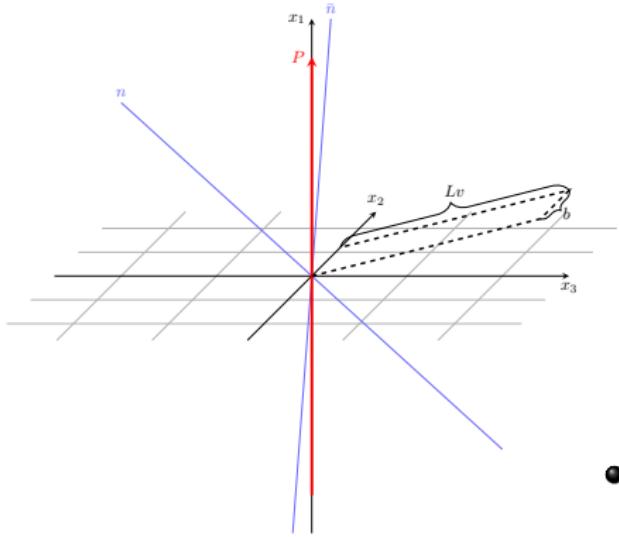
- Nucleon Matrix Element which is sensitive to transverse momentum

$$W_f^{[\Gamma]}(b; L, v; P, S) = \frac{1}{2} \langle P, S | \bar{q}_f(b) \Gamma [b, b + Lv] [b + Lv, Lv] [Lv, 0] q(0) | P, S \rangle$$

with nucleon state  $|P, S\rangle$

- Time equal matrix element  
 $v^0 = b^0 = 0$
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Musch, et. al., 1011.1213

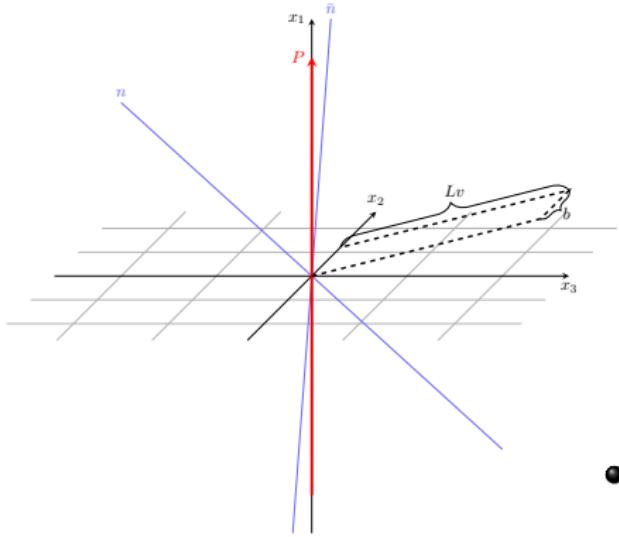


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- For our analysis we require  $b^+ = 0$  and  $\mathbf{P}_T = \mathbf{0}$  thus

$$\begin{aligned} S &= (0, 0, 0, 1), & P &= (P_0, P_1, 0, 0) \\ b &= (0, 0, b_2, b_3), & v &= (0, v_1, v_2, v_3) \end{aligned}$$

- To improve the lattice signal, HYP- and momentum smearing is applied
- u-d Quark channel to cancel disconnected diagrams
- We use the CLS ensembles with  $N_f = 2 + 1$  flavors and Wilson-Clover fermions

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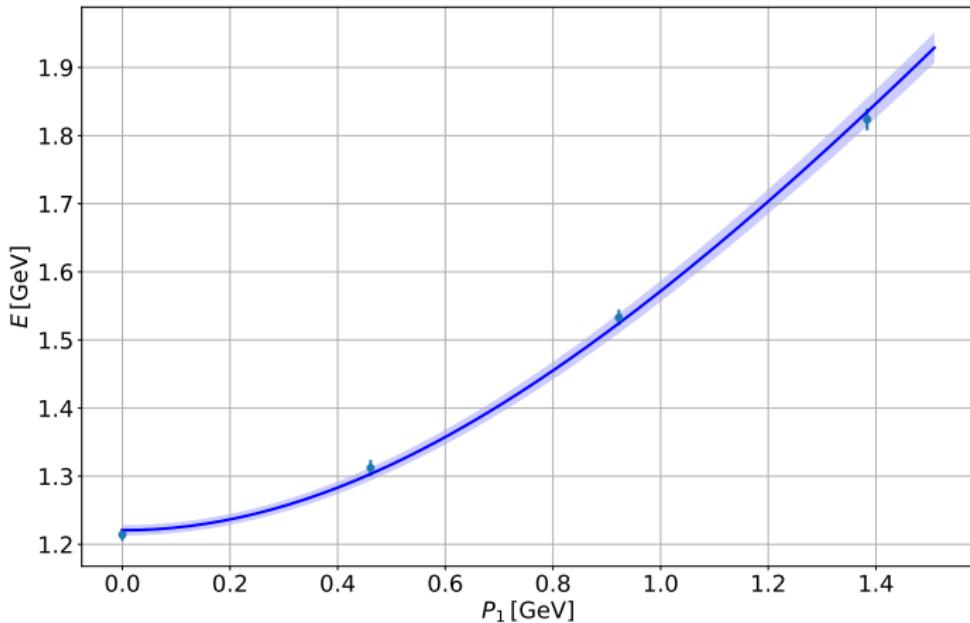
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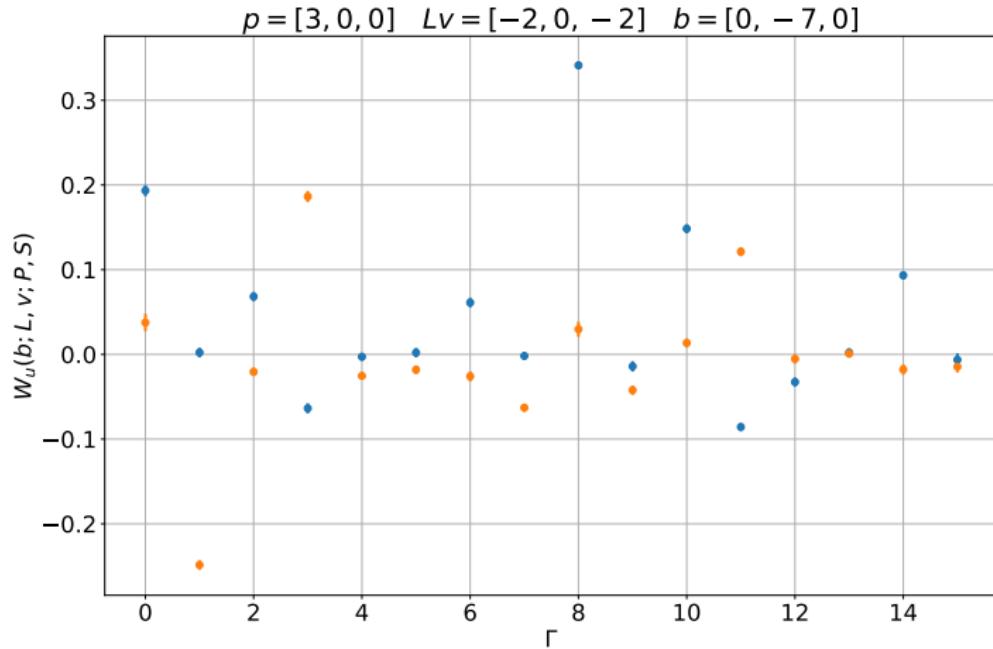
## CLS ensemble H101

$\beta$	$L^3 \times T$	$a$	$m_\pi$	#cnfg
3.4	$32^3 \times 96$	0.084fm	422MeV	2000



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The correlators are grouped into Lorentz-invariant products of  $b$ ,  $P$ ,  $Lv$ :  $(P^2, b^2, (Lv)^2, Lv \cdot P)$  and parameterized, e.g. for the vector channel  $\gamma^\mu$  as

$$\begin{aligned} \frac{1}{2}W_f^{[\gamma^\mu]}(b; L, v; P, S) = & P^\mu A_2 - im_N^2 b^\mu A_3 \\ & - im_N \epsilon^{\mu\nu\alpha\beta} P_\nu b_\alpha S_\beta A_{12} + m^2 Lv^\mu B_1 + m_N L \epsilon^{\mu\nu\alpha\beta} P_\nu v_\alpha S_\beta B_7 \\ & - im_N^3 L \epsilon^{\mu\nu\alpha\beta} b_\nu v_\alpha S_\beta B_8 - m_N^3 L (b \cdot S) \epsilon^{\mu\nu\alpha\beta} P_\nu b_\alpha v_\beta B_9 \\ & - im_N^3 L^2 (v \cdot S) \epsilon^{\mu\nu\alpha\beta} P_\nu b_\alpha v_\beta B_{10} \end{aligned}$$

Musch, et. al., 1111.4249

All  $A_i$  and  $B_i$  and  $S$  depend on  $(P^2, b^2, (Lv)^2, Lv \cdot P)$

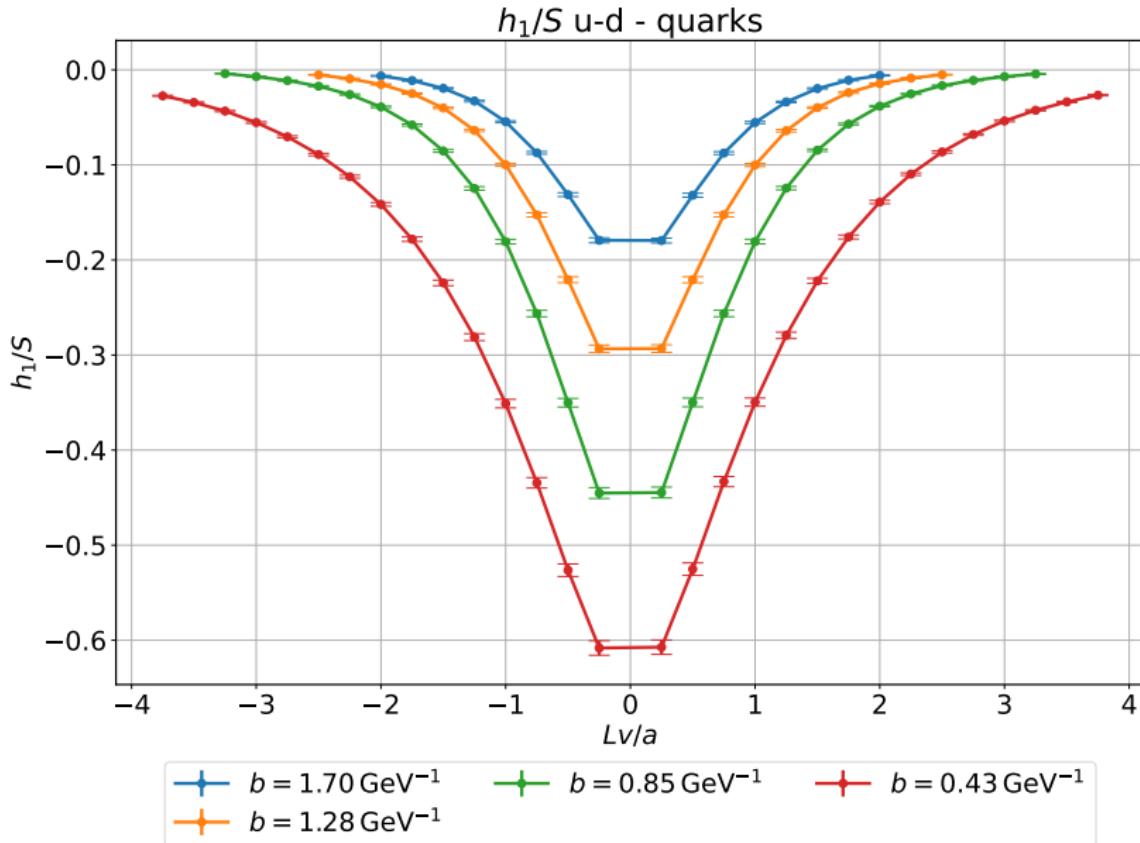
$$f_1(P^2, b^2, (Lv)^2, Lv \cdot P) = 2 \frac{A_2 + RB_1}{S(b^2, (Lv)^2, Lv \cdot P)}$$

$$g_{1L}(P^2, b^2, (Lv)^2, Lv \cdot P) = 2 \frac{A_7 + RB_{13}}{S(b^2, (Lv)^2, Lv \cdot P)}$$

$$h_1(P^2, b^2, (Lv)^2, Lv \cdot P) = 2 \frac{A_9 + RB_{15} - \frac{1}{2}m_N^2 b^2 (A_{11} - RB_{17})}{S(b^2, (Lv)^2, v \cdot P)}$$

$$R(P^2, Lv \cdot P) = \frac{m^2 Lv^+}{P^+}$$

Musch, et. al., 1111.4249



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$$R^{[\Gamma]}(Lv, b) = \left( \frac{P_2^+}{P_1^+} \right)^{-K(b,\mu)} r^{[\Gamma]} + \mathcal{O}(\lambda)$$

$$= \frac{W_f^{[\Gamma]}(P_1)}{W_f^{[\Gamma]}(P_2)}$$

where  $r^{[\Gamma]}$  is given by

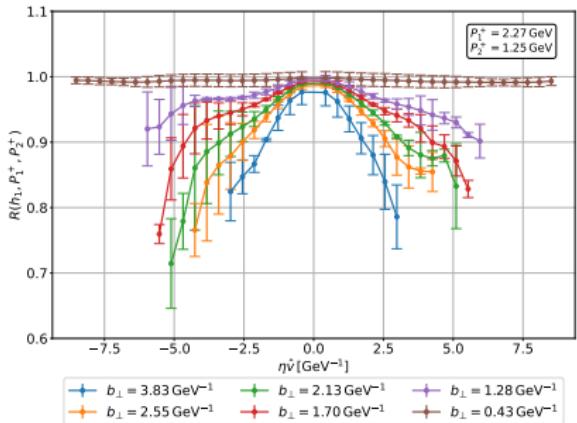
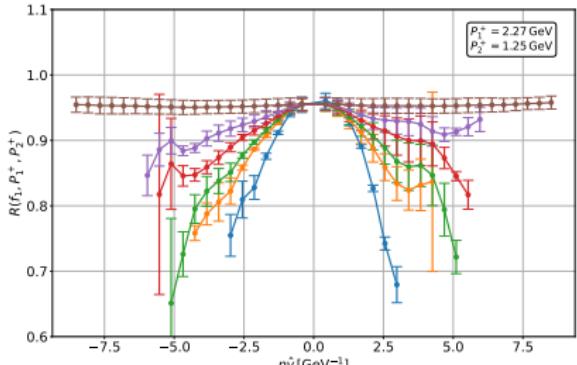
$$r^{[\Gamma]} = 1 + 4C_F \frac{\alpha_s(\mu)}{4\pi} \ln \left( \frac{P_1^+}{P_2^+} \right)$$

$$\left[ 1 - \ln \left( \frac{4P_1^+ P_2^+ |v^-|^2}{\mu^2} \right) - 2\mathbf{M}_{\ln|x|}^{[\Gamma]}(b, \mu) \right] + \mathcal{O}(\alpha_s^2)$$

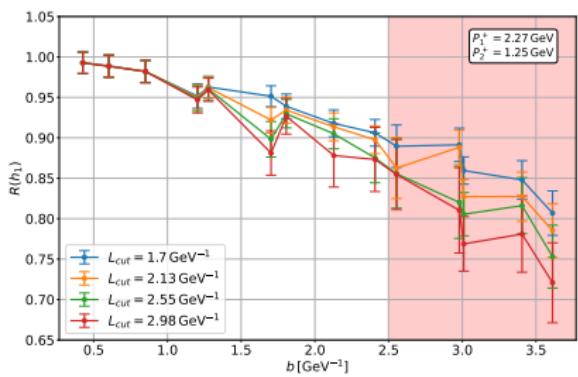
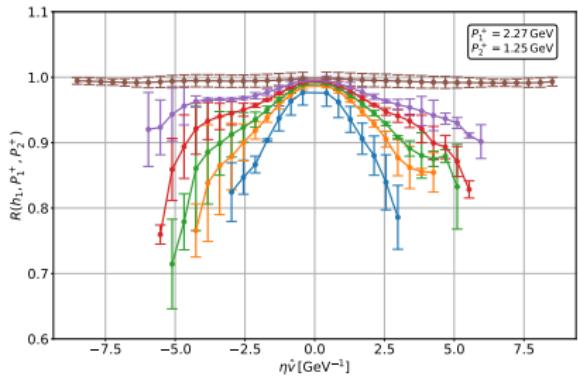
with non-perturbative  $\mathbf{M}_{\ln|x|}^{[\Gamma]}(b, \mu)$

Vladimirov, Schäfer, 2002.07527

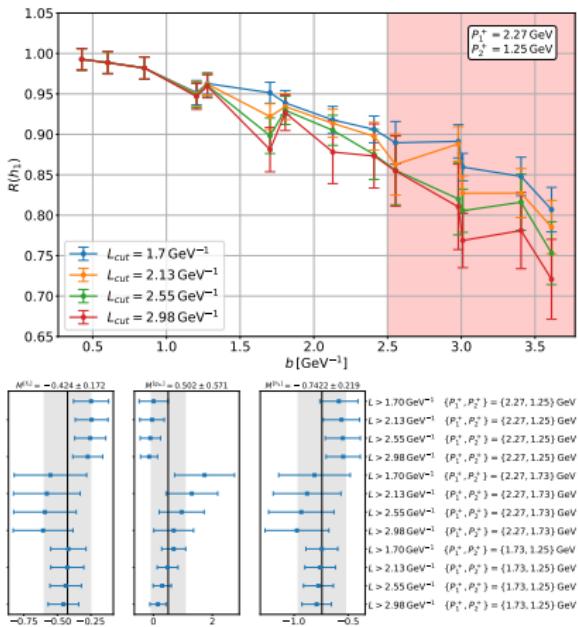
- Form ratios of  $f_1$ ,  $h_1$  and  $g_{1L}$  at different  $P^+$
- Extrapolate  $L \rightarrow \infty$
- Use the lattice data to determine  $M_{\ln|x|}^{(n),\Gamma}(b, \mu)$  at approx  $1\text{GeV}^{-1}$

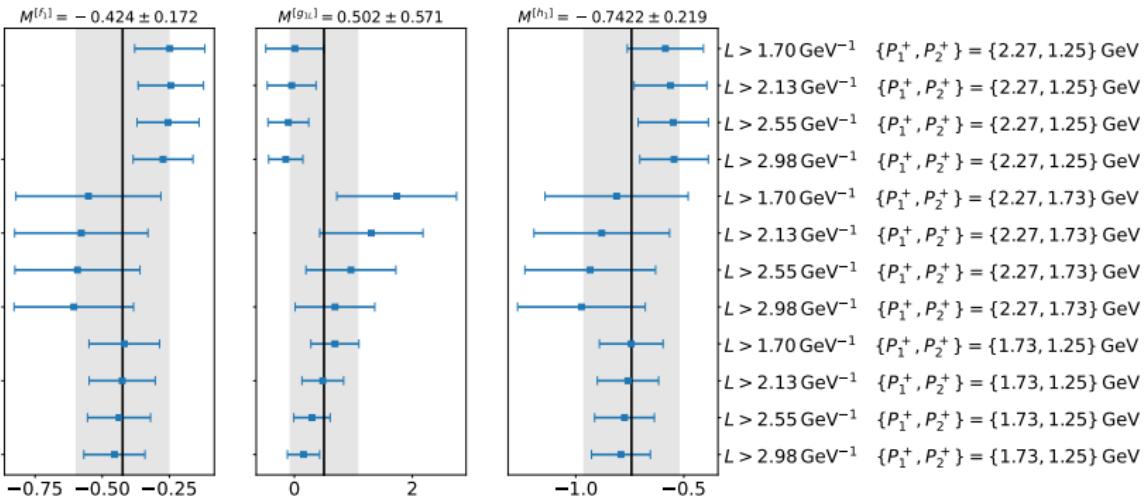


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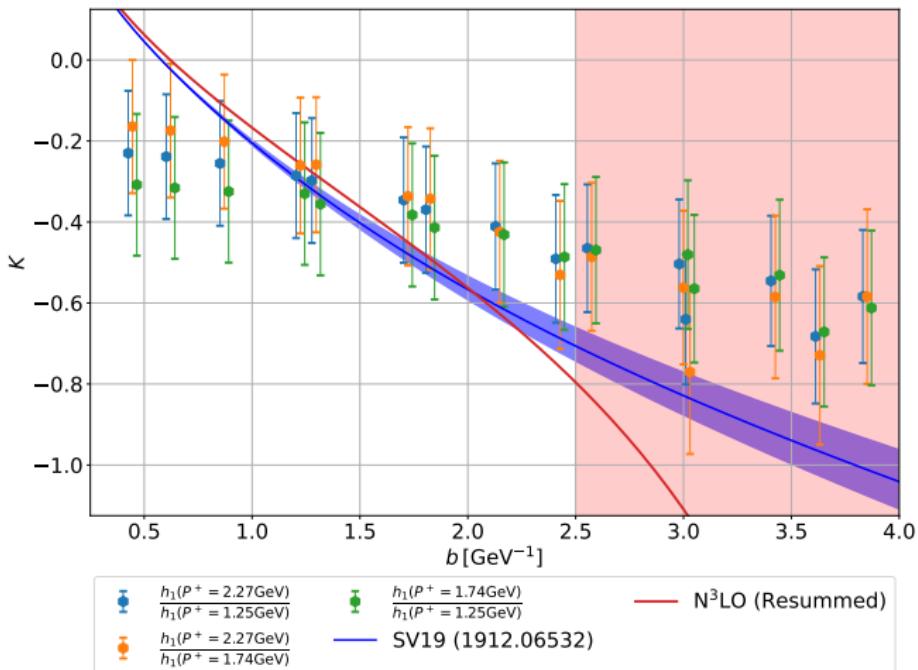
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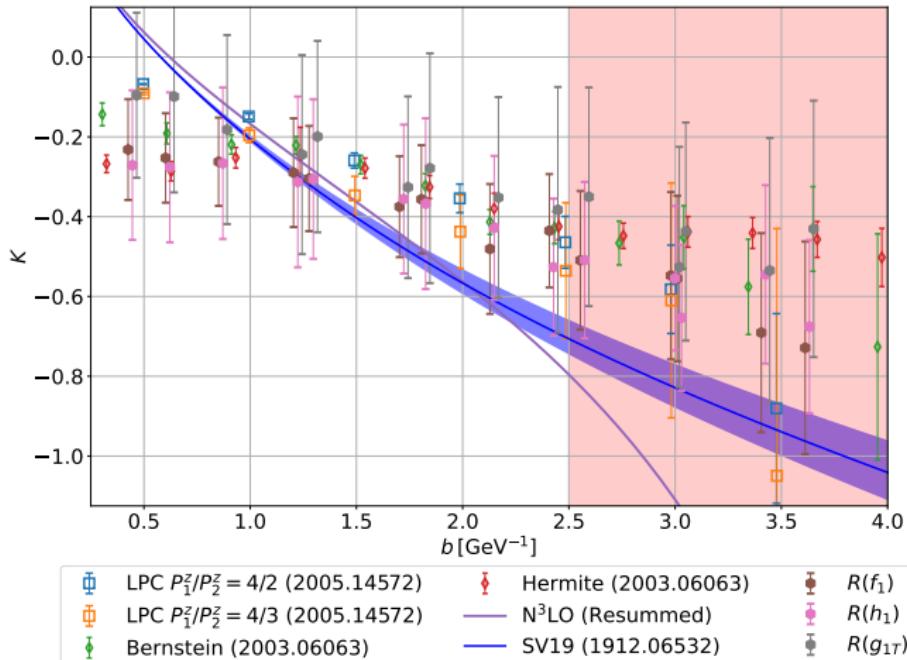
Use  $\mathbf{M}_{\ln|x|}^{[\Gamma]}(b, \mu)$  to compute  $r^{[\Gamma]}$   
and extract  $K$

$$K = \frac{\ln \left( \frac{r^{[\Gamma]}}{R^{[\Gamma]}} \right)}{\ln \left( \frac{P_2^+}{P_1^+} \right)}$$



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- Systematic errors need to be reduced, i.e. better understood
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Thank you for your attention  
Questions?

